

# Berry Phase & Chern Numbers

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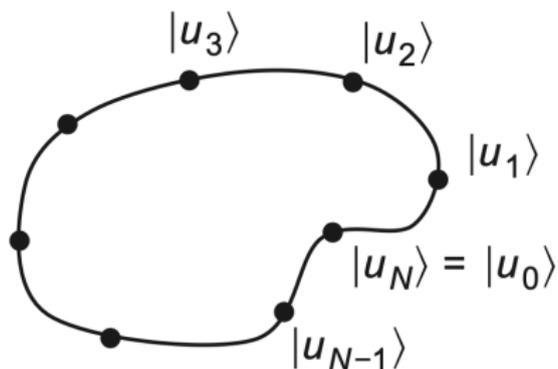
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- Berry Phase is the total amount of phase accumulated by the loop.

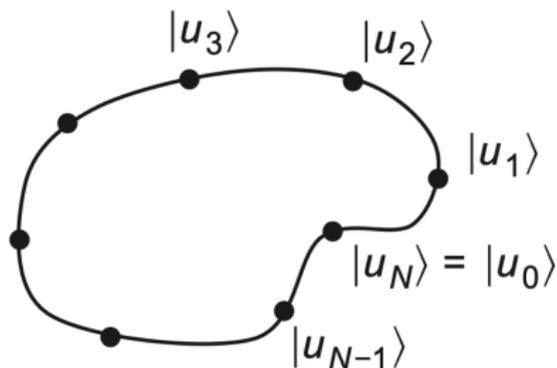
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## Berry Phase

We may define the Berry Phase as:

$$BP := -\arg \left[ \langle u_0 | u_1 \rangle \langle u_1 | u_2 \rangle \cdots \langle u_{N-1} | u_0 \rangle \right] \quad (1)$$

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- The numbers will be different, yet nothing changed about your \*driving experience\*.

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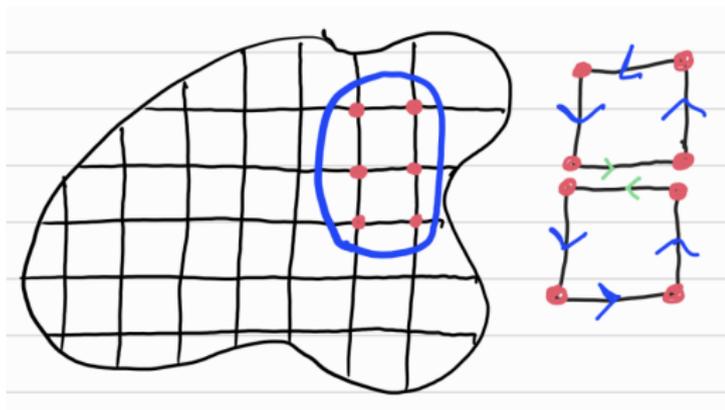
## B.P. Gauge Independence

Applying the gauge transformation to every state:

$$\begin{aligned}BP &:= -\arg \left[ \langle u_0 | u_1 \rangle \langle u_1 | u_2 \rangle \cdots \langle u_{n-1} | u_0 \rangle \right] \\ &\mapsto -\arg \left[ \left( e^{i(-\beta_0 + \beta_1 - \beta_1 + \beta_2 - \beta_2 + \cdots + \beta_0)} \langle u_0 | u_1 \rangle \langle u_1 | u_2 \rangle \cdots \langle u_{n-1} | u_0 \rangle \right) \right] \\ &= -\arg \left[ e^0 \langle u_0 | u_1 \rangle \langle u_1 | u_2 \rangle \cdots \langle u_{n-1} | u_0 \rangle \right] \\ &= BP\end{aligned}$$

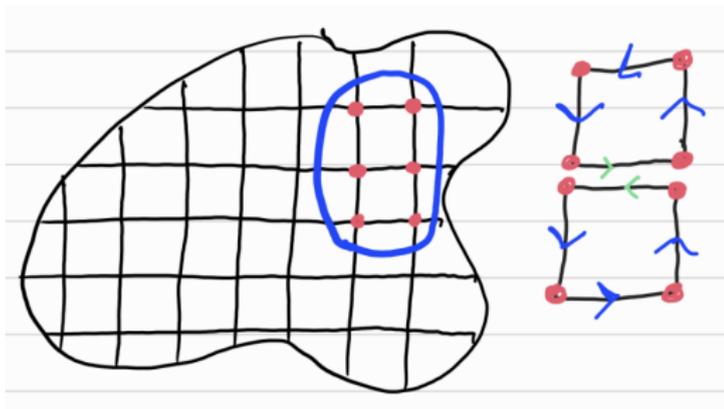
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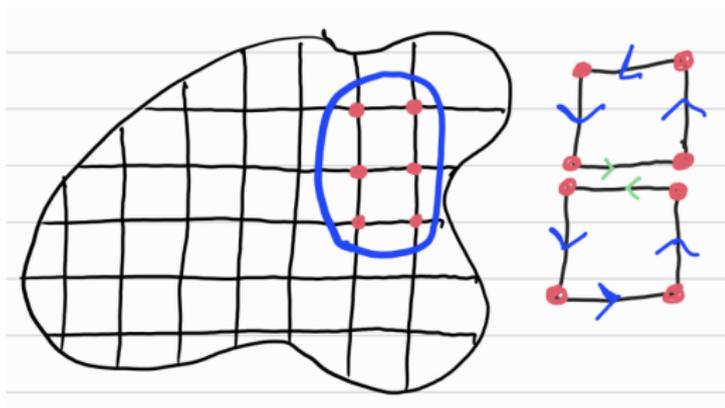
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- Let's calculate the Berry Phase of the loop along the boundary.
- Alt. Approach: Add up BP of the enclosed plaquettes.
  - Each internal edges is included twice with opposite orientations.
  - The BP for a plaquette,  $\square$ , is called *Berry Flux*, denoted as  $F_{\square}$ .

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- This number  $Q$  is the Chern Number.

# Continuous Formulation

- $\ln \langle u_\lambda | u_{\lambda+d\lambda} \rangle = \ln \langle u_\lambda | (|u_\lambda\rangle + d\lambda |\partial_\lambda u_\lambda\rangle + \dots) \rangle \approx d\lambda \langle u_\lambda | \partial_\lambda u_\lambda \rangle$ 
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- The integrand is purely imaginary  $\implies \gamma_L = \oint_L i \langle u_\lambda | \partial_\lambda u_\lambda \rangle d\lambda$

$$\begin{aligned} 2\text{Re} \langle u_\lambda | \partial_\lambda u_\lambda \rangle &= \langle u_\lambda | \partial_\lambda u_\lambda \rangle + \overline{\langle u_\lambda | \partial_\lambda u_\lambda \rangle} \\ &= \langle u_\lambda | \partial_\lambda u_\lambda \rangle + \langle \partial_\lambda u_\lambda | u_\lambda \rangle \\ &= \partial_\lambda \langle u_\lambda | u_\lambda \rangle = 0 \end{aligned}$$

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- For  $\dim \geq 2$  define the *Berry Curvature* via the curl:

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- $\Omega(\lambda)$  is Gauge Invariant
- Apply Stokes' Theorem (to what you can):

$$\oint_L \vec{A}(\lambda) \cdot d\lambda + 2\pi Q = \iint_S \Omega(\lambda) dS$$

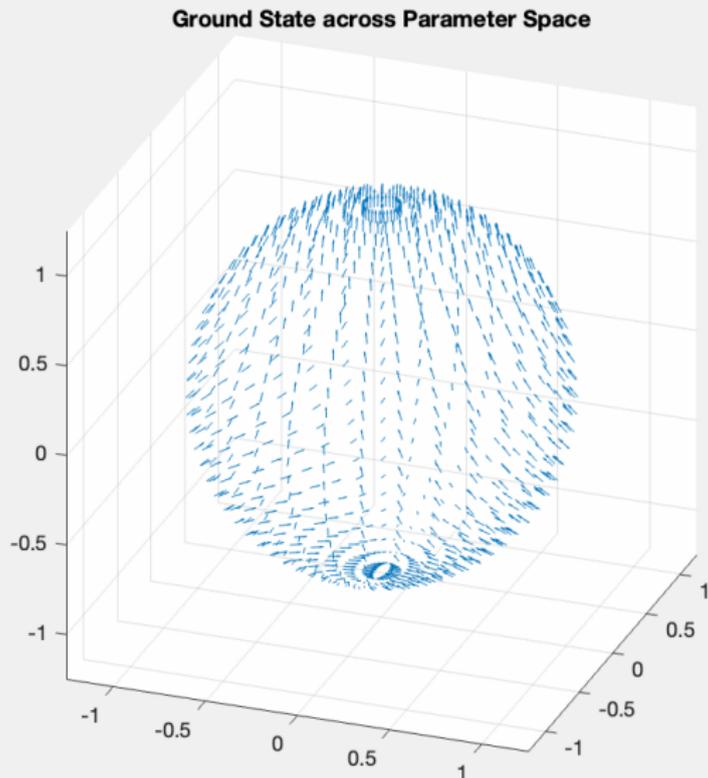
## Chern Theorem

For a closed 2D manifold we have that the integral of the Berry Curvature is  $2\pi Q$ , where  $Q \in \mathbb{Z}$  is the Chern Number.

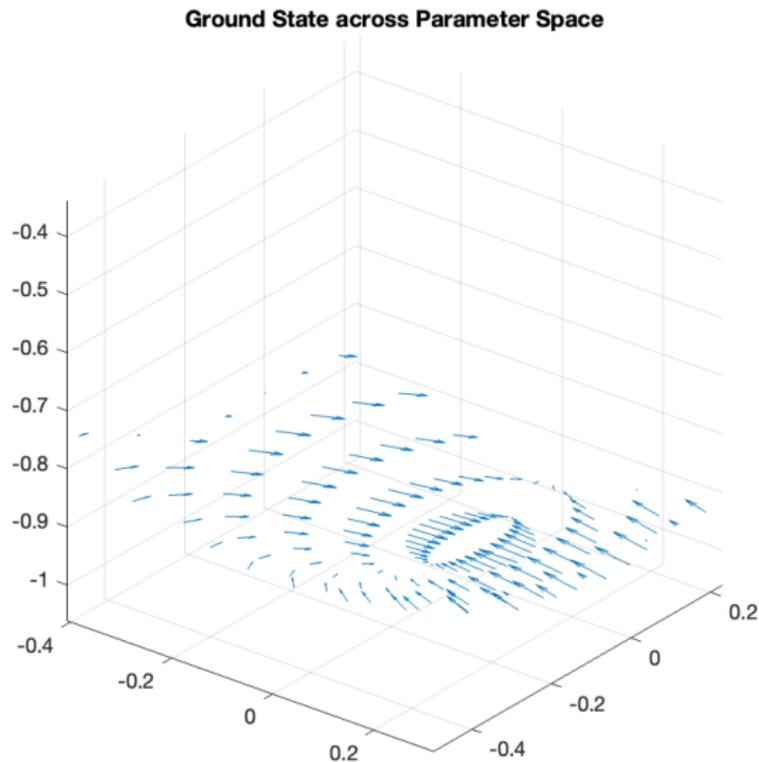
$$\iint_S \Omega(\lambda) \cdot d\mathbf{S} = 2\pi Q \quad (3)$$

"Note that when the Chern number is nonzero, it is impossible to construct a smooth and continuous gauge over the entire surface  $S$ " [Vanderbilt].

# Gauge Obstruction (1)

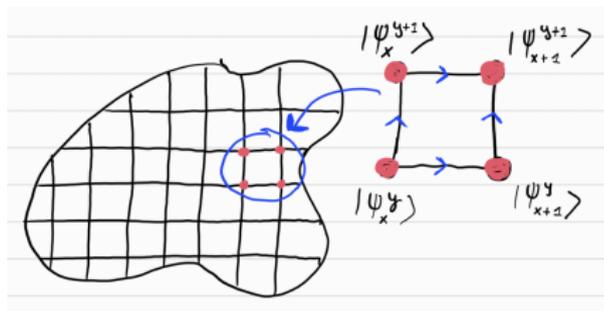


# Gauge Obstruction (2)



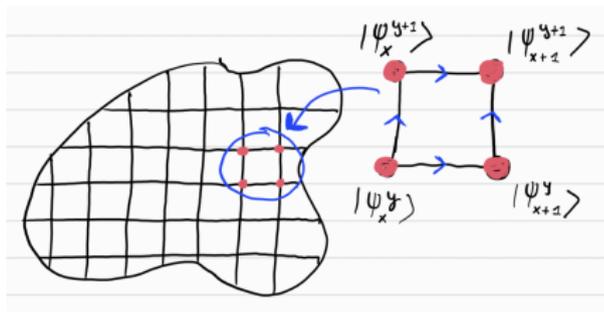
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- Step 1: Discretize the Torus into plaquettes.



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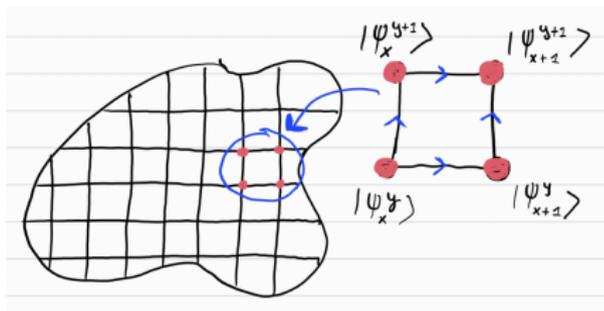
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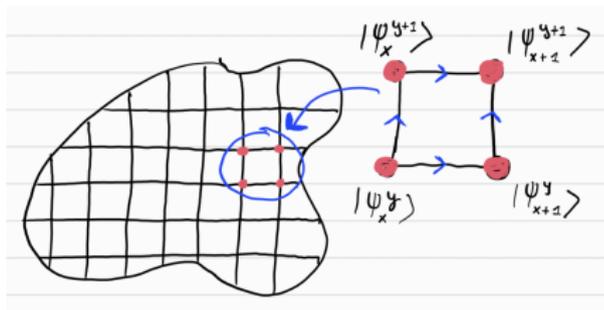


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- Step 4: Calculate Berry Flux for each plaquette and sum:

$$\tilde{F}_x^y = \ln \left( U(e_1)U(e_2)U(e_3)^{-1}U(e_4)^{-1} \right) \quad (4)$$

# Application (1)

- We follow Fukui et al. and apply the method to the following family of Hamiltonians parameterized by  $\vec{\mathbf{k}} = (k_x, k_y)$ :

$$H(\vec{\mathbf{k}}) = \begin{pmatrix} -2t \cos(k_y - \frac{2}{3}\pi) & -t & -te^{-3ik_x} \\ -t & -2t \cos(k_y - \frac{4}{3}\pi) & -t \\ -te^{3ik_x} & -t & -2t \cos(k_y - 2\pi) \end{pmatrix}$$

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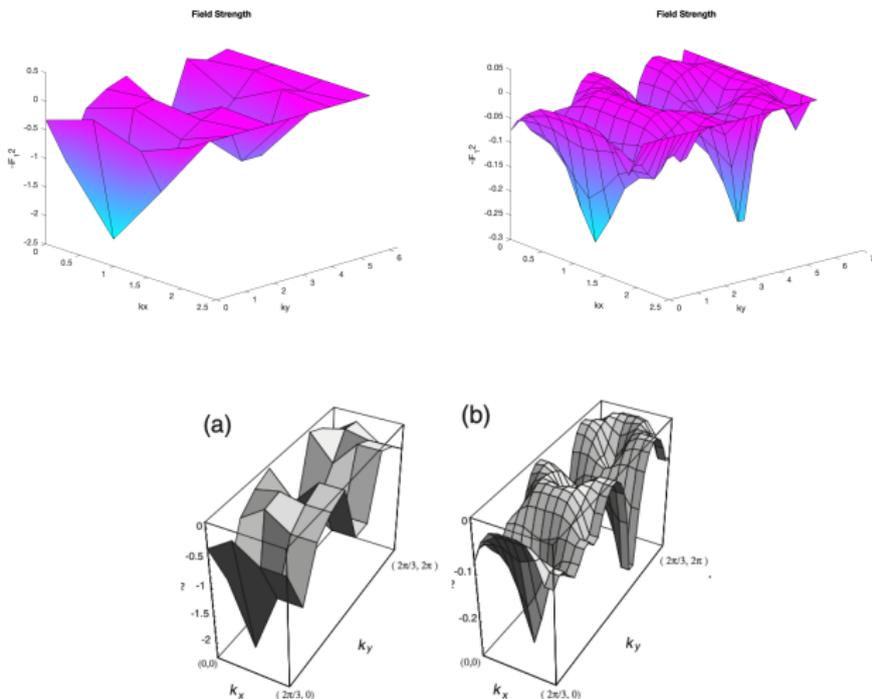
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- This family of matrices is obtained from the floquet representation of the Hamiltonian describing spinless fermions subjected to an eternal magnetic field undergoing certain flux constraints.

# Results

We reproduced the results of Fukui et al., i.e.  $Q = -2$  and found similar  $\tilde{F}$ -surfaces:



## Application (2)

- Qi-Wu-Zhang Model is given by a Hamiltonian in  $k$ -space with an additional parameter,  $u$ :

$$H(\vec{\mathbf{k}}, u) = \begin{pmatrix} u + \cos(k_x) + \cos(k_y) & \sin(k_x) - i \sin(k_y) \\ \sin(k_x) + i \sin(k_y) & -u - \cos(k_x) - \cos(k_y) \end{pmatrix}$$

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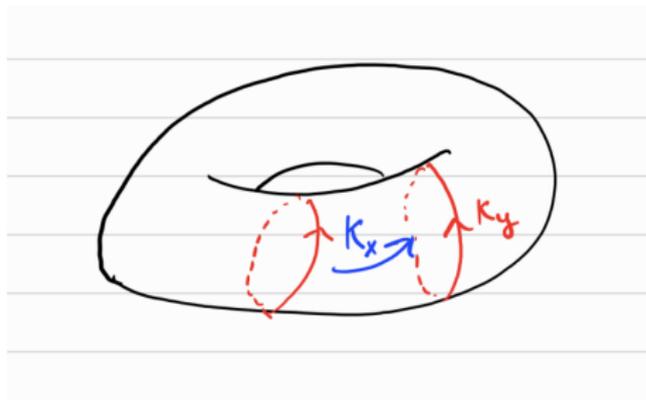
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- The Chern Number as a function of  $u$  for this model is:

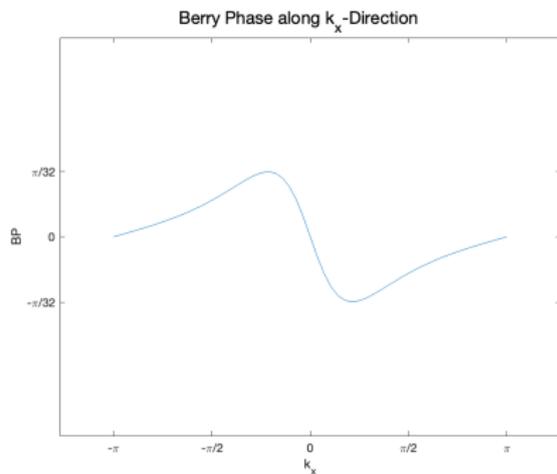
$$Q(u) = \begin{cases} 0 & \text{if } |u| > 2 \\ -1 & \text{if } -2 < u < 0 \\ +1 & \text{if } 0 < u < 2 \end{cases}$$

Which is exactly what we see in our code.

# Visualizing QWZ Model's Chern Number



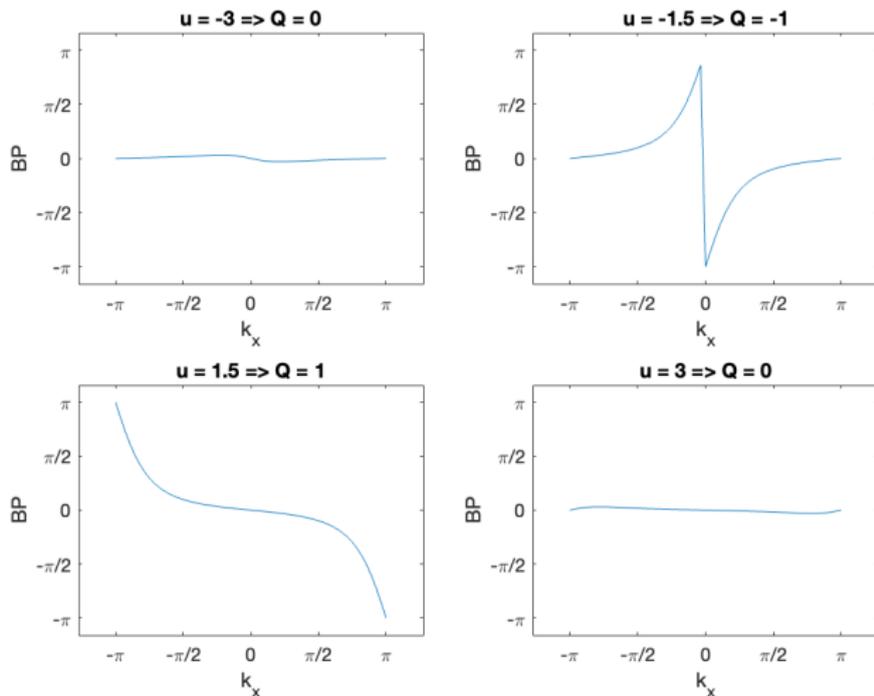
(a) Brillouin Zone



(b) BP along rings of the torus.

# Visualizing QWZ Model's Chern Number

## Berry Phase in $k_x$ -Direction



# References I

-  Michael Berry, *Quantal Phase factors for accompanying adiabatic changes*, Royal Society, 1984.
-  Takahiro Fukui, Yasuhiro Hatsugai, and Hiroshi Suzuki, *Chern Numbers in Discretized Brillouin Zone: Efficient Method of Computing (Spin) Hall Conductances*, Journal of the Physical Society of Japan Vol. 74, No. 6, June, 2005.
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